

Ice dynamics and basal properties of Sofiyskiy glacier, Altai mountains, Russia, based on DGPS and radio-echo sounding surveys

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ABSTRACT. Central Asian glaciers, compared to other glaciers in the world, should exhibit a different response to changing climate owing to (i) the extreme continentality that gives rise to aridity and large seasonal temperature variations, and (ii) the coincidence of accumulation and ablation seasons in summer. A detailed 4 year field survey has been carried out on Sofiyskiy glacier in the south Chuya range (Altai mountains), Russia. Field observations reveal that this glacier retreated steadily during the 20th century at a rate of 18.3 m a^{-1} , considerably faster than Maliy Aktru glacier, 30 km to the north. Based on radio-echo sounding (RES), the subglacial topography of Sofiyskiy glacier has been reconstructed, and stake and snow-pit measurements enable the determination of surface velocity and mass balance. Basal sliding plays an important role in Sofiyskiy glacier's behaviour, as is shown by a force-balance analysis. An analysis of RES measurements yields the bed reflection power (BRP). Using the three-layer reflectivity model of Born and Wolf (1986), the BRP is compared with the theoretical bed reflectivity for two contrasting layers of varying porosity and thickness. Results indicate that Sofiyskiy glacier possibly exhibits a polythermal regime.

INTRODUCTION

Glaciers are potentially good climatic indicators. However, most of the source information on the dynamic response of temperate glaciers to changing climate comes from European and, to a lesser extent, North American glaciers, although European glaciers form only 3% of the world glacier area (Dyrgerov and Meier, 1997). Hence, our knowledge of the glaciers of remote and inaccessible mountain regions, often characterized by a high degree of glacierization, is rather limited. This biased distribution of glaciological source information severely constrains present attempts to assess the impact of climate change on the global glacier volume. It also prevents a full understanding of glacier dynamics and the impact of different climatic conditions. This paper seeks to expand the current global glaciological record by presenting and analyzing data gathered on Sofiyskiy glacier, located in the Russian Altai mountains (Fig. 1), during four consecutive summer field seasons in the period 1997–2000. The climate of the Altai mountains has two important characteristics which make this study area especially interesting. First, the extreme continentality gives rise to aridity and a large diurnal and seasonal temperature range. Second, owing to the coincidence of accumulation and ablation seasons in summer, we can characterize the Altai glaciers as *summer-accumulation*-type glaciers. Although current knowledge of the behaviour of these glaciers is limited, they are reported

to be more vulnerable to the current global warming trend than the better-known (maritime) *winter-accumulation*-type glaciers (Fujita and Ageta, 2000). In this paper, the basic characteristics of Sofiyskiy glacier are discussed, such as surface mass balance and velocity, and the historical front variations. To improve understanding of the glacier's dynamics,

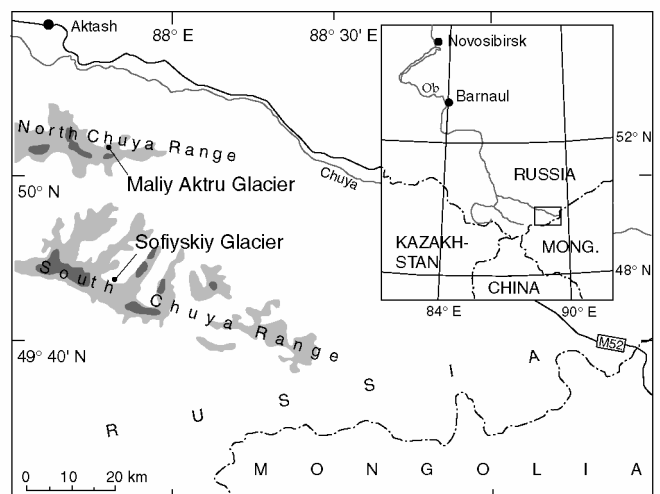


Fig. 1. Location map of the Altai mountains and the study area (Sofiyskiy glacier).

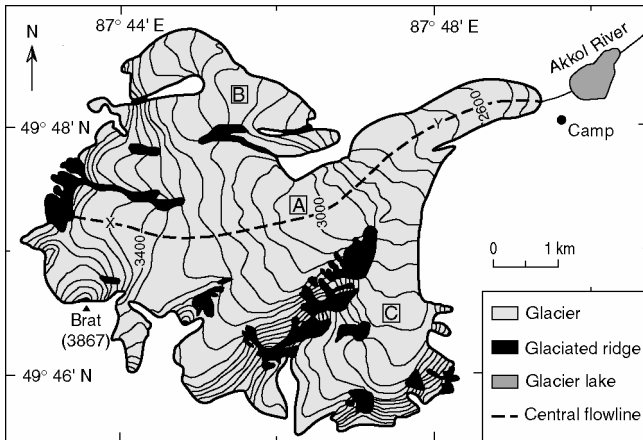


Fig. 2. Contour map of Sofiyskiy glacier, based on the 1952 topographic map. At present, basins B and C are no longer connected with the main ice flow.

we also present a force-balance analysis and an analysis of basal conditions using a three-layer reflectivity model for electromagnetic pulses.

STUDY AREA AND RECENT GLACIAL HISTORY

Sofiyskiy glacier (49°47' N, 87°46' E) is one of the major glaciers in the central part of the south Chuya range, a mountain chain within the Altai mountain belt, situated in the border region between Russia and Mongolia (Fig. 1). At present, the glacierized area consists of three basins (Fig. 2), i.e. the main basin (A) and two smaller basins (B and C). The latter are at present no longer connected to the main ice flow. The southwest–northeast-oriented Sofiyskiy glacier is approximately 7 km long and has an average width of 1.7 km. The head of the glacier is located near Brat mountain (3876 m a.s.l.), one of the highest peaks in the range. The cumulation area has a mean surface slope of $\sim 17\%$ at the central flowline. A steep icefall is situated between 2800 and 3000 m a.s.l. Below this icefall, the glacier terminates in a relatively flat tongue.

During the 20th century, several front-position markers were recorded. The earliest records are from field markers and aerial photographs (Bulatov and others, 1967). During each of the 1997–2000 field seasons, the snout position was

Table 1. Historic front positions of Sofiyskiy glacier, measured as distance from present front position

Position m	Time AD	Source
9395	-680 ± 60	^{14}C dating
3130	1240 ± 35	^{14}C dating
2561	1630 ± 60	^{14}C dating
1871	1898	Field marker
734	1939	Field marker
553	1952	Field marker
318	1963	Field marker
61	1997	DGPS survey
40	1998	DGPS survey
7	1999	DGPS survey
0	2000	DGPS survey

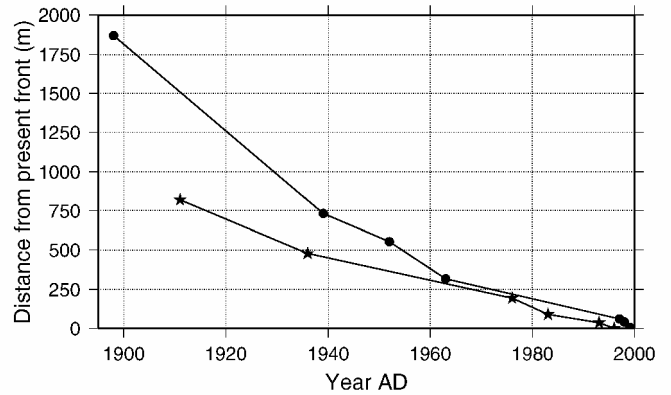


Fig. 3. Historic front positions of Sofiyskiy glacier (●) and Maliy Aktru glacier (★) during the 20th century.

mapped with differential global positioning system (DGPS). The historical record is extended further back in time by radiocarbon dating of wood fragments collected in three different former front moraines of Sofiyskiy glacier (Table 1). Comparison with Maliy Aktru glacier, located in the north Chuya range some 30 km to the northwest (Fig. 1), shows that both glaciers have retreated steadily since the beginning of the 20th century. Average retreat rates are 9.6 m a^{-1} for Maliy Aktru glacier and 18.3 m a^{-1} for Sofiyskiy glacier, based on remeasurement of field markers and our own measurements of the front position. The most rapid retreat of both glaciers occurred between 1900 and 1940 (Fig. 3).

SURFACE MASS BALANCE AND GLACIER DYNAMICS

A network of about 35 stakes was employed for determining local mass balance and ice velocity. Owing to the presence of many hidden crevasses and severe snow conditions in the accumulation area, this network was confined to the zone below 3100 m a.s.l. The position of the stakes was measured with a Leica SR 9500 DGPS. This system consists of a fixed reference station (“Camp” in Fig. 2) and a remote “rover” station. The system assures a centimeter accuracy in hori-

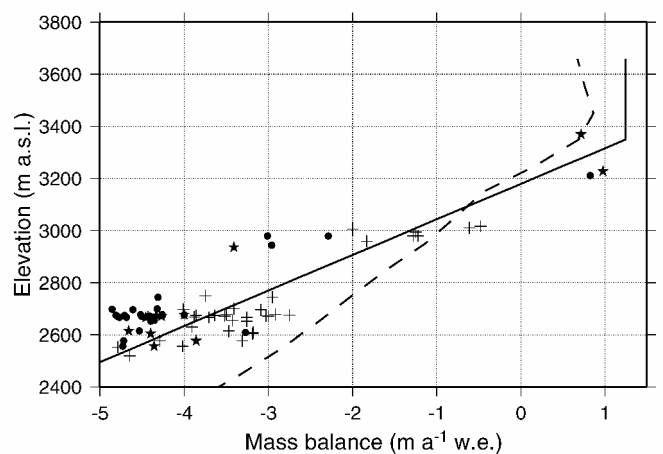


Fig. 4. Mass-balance measurements carried out on Sofiyskiy glacier for the mass-balance years 1997/98 (★), 1998/99 (●) and 1999/2000 (+). The solid line is the best fit of these data. The dashed line is the mean mass-balance profile for Maliy Aktru glacier based on a 38 year mass-balance record (Hoelzle and Haerberli, 1999).

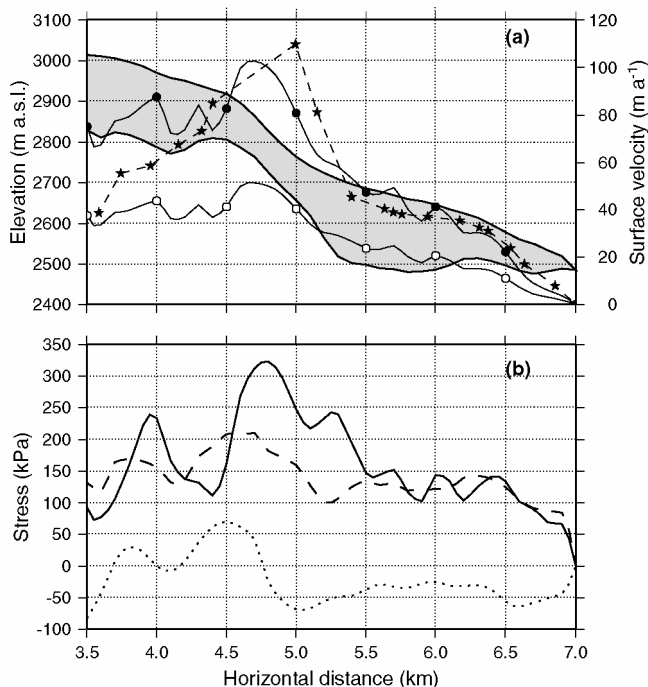


Fig. 5. (a) Surface and bedrock profile (gray shaded), surface velocity measurements (★) and modeled surface velocities using $A = 0.8 \times 10^{-16} \text{ Pa}^{-3} \text{ a}^{-1}$ (white circles) and using $A = 1.6 \times 10^{-16} \text{ Pa}^{-3} \text{ a}^{-1}$ (black circles) along the central flowline of Sofiyskiy glacier. (b) Driving stress τ_d (solid line), predicted basal shear stress τ_{xz} (dashed line) and predicted surface longitudinal stress deviator τ'_{xx} calculated with the model along the central flowline.

zonal position and an accuracy of 3–5 cm in elevation. Base lengths were always <10 km and usually <5 km, further limiting the error. However, taking into account errors due to the repeat survey, and positioning errors of the rover station due to surface melt, the overall accuracy should be somewhat lower.

Glacier mass-balance measurements are conventionally carried out at the end of the ablation season. Due to logistic restrictions, these mass-balance measurements had to be carried out some weeks beforehand. Accumulation was determined from a number of snow pits dug in 1999. Figure 4 shows the individual mass-balance measurements for each balance year. It should be noted that the ablation values for 1999–2000 are probably underestimated, as the ablation stakes were remeasured early in the summer season. Based on the individual measurements, a mean mass-balance reference profile for the period 1997–2000 was constructed using a linear parameterization (Fig. 4). Excluding the short 1999/2000 balance season, the equilibrium-line altitude (ELA) and the slope of this reference profile are 3179 m and $73.1 \text{ cm a}^{-1} \text{ w.e. per } 100 \text{ m}$, respectively. Figure 4 also shows the mass-balance reference profile for Maliy Aktru glacier, based on a 38 year mass-balance record (1962–99; Hoelzle and Haeberli, 1999). The mean ablation at the glacier tongue is distinctly lower for Maliy Aktru than for Sofiyskiy glacier. This might be due to the fact that Maliy Aktru glacier lies in a very narrow and steep valley, and is therefore more shaded from direct radiation.

Based on these stake measurements, a profile of the average horizontal surface velocities along the central flowline was compiled (Fig. 5). This profile shows a maximum

surface velocity of 110 m a^{-1} at the icefall; the average surface velocity is 35 m a^{-1} over the glacier tongue and 55 m a^{-1} over the entire ofile.

FORCE-BALANCE ANALYSIS

To gain a better insight into Sofiyskiy glacier dynamics we applied the higher-order numerical flowline model of Pattyn (2002) in a diagnostic way. Besides shear stresses, this model takes into account longitudinal stress gradients in calculating the ice deformation, and uses Glen's flow law as a constitutive equation. Based on the local ice geometry along the central flowline (surface and bedrock topography, width of the drainage basin), a constant flow parameter (experimentally determined as $A = 0.8 \times 10^{-16} \text{ Pa}^{-3} \text{ a}^{-1}$; De Smedt and Pattyn, 2003), and considering a no-slip basal boundary condition, the full stress field was determined. Major components of the stress field are displayed in Figure 5. The icefall area is characterized by high driving stresses ($\tau_d > 300 \text{ kPa}$), and for most of the long profile $\tau_d > 100 \text{ kPa}$. The basal shear stress is a smoothed version of the driving stress, characterized not only by a lower amplitude, but also by phase shifts compared to the driving-stress profile. Such phase shifts are due to local pushes and pulls as a result of bedrock and surface perturbations (Pattyn, 2002). Longitudinal stretching is confined to the region upstream of the icefall. The downstream area (ice tongue) is dominated by compressive flow.

Modeled surface velocities are in agreement with the measured values upstream of the icefall. For the whole downstream section, measured surface velocities are significantly higher, especially in the icefall region. The discrepancy is probably due to neglect of basal sliding in the model (no-slip boundary condition). We carried out an experiment with a higher value of $A = 1.6 \times 10^{-16} \text{ Pa}^{-3} \text{ a}^{-1}$, hence making the ice softer (warmer). While softer ice accounts for a better agreement between modeled and observed surface velocities in the ice tongue, the discrepancy becomes larger upstream of the icefall (modeled velocity too high) and at the icefall itself (modeled velocity still too low). These results thus indicate that (i) deformational velocity alone cannot account for the spatial distribution of the velocity field of the glacier and that a significant amount of the total surface velocity must be due to basal sliding in the downstream area of the glacier, or (ii) the glacier ice is softer (hence warmer) in the ablation area compared to the area upstream of the icefall.

RADIO-ECHO SOUNDING

Radio-echo sounding (RES) measurements were carried out with a 5 MHz (central frequency) ice-penetrating radar (Narod and Clarke, 1994). The system consists of a mono-pulse transmitter generating 1600 V pulses across a resistively loaded 10 m dipole antenna, and an airwave-triggered oscilloscope receiver connected to a palmtop computer for digital data recording. Pulses are generated at 512 Hz and have a bandwidth of 1–200 MHz. Similar systems have already been used successfully on several glaciers in Canada and Alaska (Narod and Clarke, 1994; Welch and others, 1998; Copland and Sharp, 2001). Apart from a number of isolated RES measurements, four profiles were taken in crevasse-free zones, which reveal details of the bed topography and basal conditions. At $\sim 3450 \text{ m}$ (“X” in Fig. 2), a longitudinal (along the central flowline) and a transverse radar

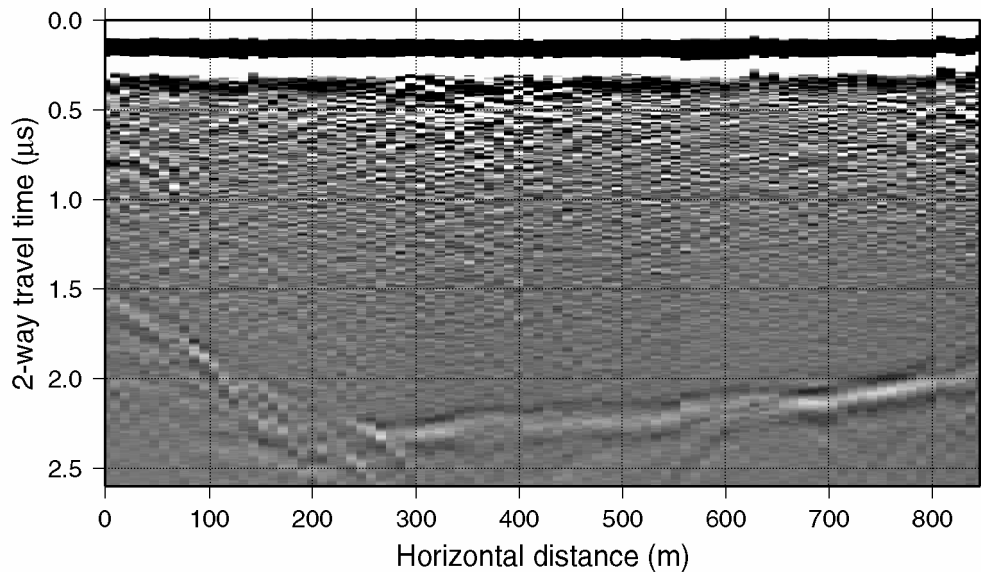


Fig. 6. Line-intensity visualization of the longitudinal RES profile in the ablation area of Sofiyyskiy glacier.

profile were recorded. The horizontal distance between successive measurements was ~ 19 m in both profiles. In the ablation zone beneath the icefall, a second longitudinal-transverse profile set was recorded with a horizontal interval of ~ 9.5 m (the central point is “Y” in Fig. 2).

On all measurements, static corrections were performed for the trigger delay caused by the 38 m separation between transmitter and receiver. Band-pass filtering was not carried out. Migration was calculated along the one-dimensional profiles, but is not considered in the analysis below due to the complex topography in some places. As outlined in Copland and Sharp (2001), migration would only serve to increase values of high bed reflection power (BRP), as it tends to increase ice thickness in zones of steep basal topography. All recorded traces present relatively clear basal reflections. An example of this is given in Figure 6, which shows a line-intensity plot for the longitudinal glacier profile in the ablation zone. The two-way travel time and amplitude of the bed reflection wave were determined using a semi-automatic bed-reflection picking routine. Assuming that the bed reflection occurs from a point directly below the midpoint between the transmitter and the receiver, ice thickness was inferred using a wave propagation velocity of $168 \mu\text{m s}^{-1}$ for glacier ice, a value confirmed by a common-midpoint experiment carried out in the ablation area. A density profile in the accumulation area was obtained from Fujii and others (2002), and a propagation velocity correction was determined by considering firn as a mixture of ice and air. According to Looyenga (1965), the dielectric permittivity for such a heterogeneous mixture is given by (e.g. Hempel and others, 2000)

$$\varepsilon_{\text{firn}} = \left[\left(\varepsilon_{\text{ice}}^{1/3} - 1 \right) \frac{\rho_{\text{firn}}}{\rho_{\text{ice}}} + 1 \right]^3, \quad (1)$$

where $\varepsilon_{\text{ice}} = 3.17$ is the dielectric permittivity of ice. Combining the local ice thickness with the DGPS registration of the surface position of the traces yields four basal profiles (Fig. 7). Owing to the large number of multiple basal reflections within the reflected RES signals (most likely caused by a complex three-dimensional topography) and the relatively large distance between two successive measurements

(~ 19 m), the accuracy of the bed position in the accumulation area varies between 15 and 45 m, a value estimated from the uncertainty with which the bed position could be determined due to multiple reflectors. In areas with clear basal reflections, the uncertainty is typically of the order of one-tenth of the transmitted wavelength (Bogorodsky and others, 1985), leading to an error of ± 3.4 m. Our measurements suggest a maximum ice thickness of > 200 m in the accumulation zone (Fig. 7a). Bedrock signals are distinctly clearer in the ablation zone. Here, the 900 m longitudinal profile starts at the bottom of the icefall, where basal reflections are disturbed (Figs 6 and 7d). This is most probably due to the steepness and roughness of the local glacier bed and internal structures within the ice. Beyond the icefall, the bed steepness suddenly reduces to 4° , and basal reflections become more defined.

To gain insight into the basal layer, it is convenient to calculate the reflected power. We consider the reflected signal to be consistent with a monochromatic sinusoidal wave train (Gades and others, 2000). As such, we can define the reflected power P within a time window $t_1 - t_2$ as

$$P \equiv \frac{1}{2(t_2 - t_1 + 1)} \sum_{i=t_1}^{t_2} A_i^2, \quad (2)$$

where A_i is the recorded amplitude of the wave. Using Equation (2), we can calculate the reflection power of both the bed reflection wave (BRP) and the internal reflection waves (IRP), which are located between the end of the direct airwave and the beginning of the bed reflection wave (Gades and others, 2000). The BRP contains information on the structure and composition of the glacier bed. The IRP is a measure of the scatter and attenuation within the glacier ice. Contrary to Gades and others (2000), a variable time window is used for the calculation of the IRP, which is more suitable for shallow glaciers (Copland and Sharp, 2001). This window extends from $0.6 \mu\text{s}$ after the beginning of the airwave to $0.1 \mu\text{s}$ before the bed reflection wave. The BRP starts at the latter point and ends at $0.4 \mu\text{s}$ after the bed reflection wave.

If we assume that losses due to polarization and scattering at the ice/bed interface are minimal or constant, and

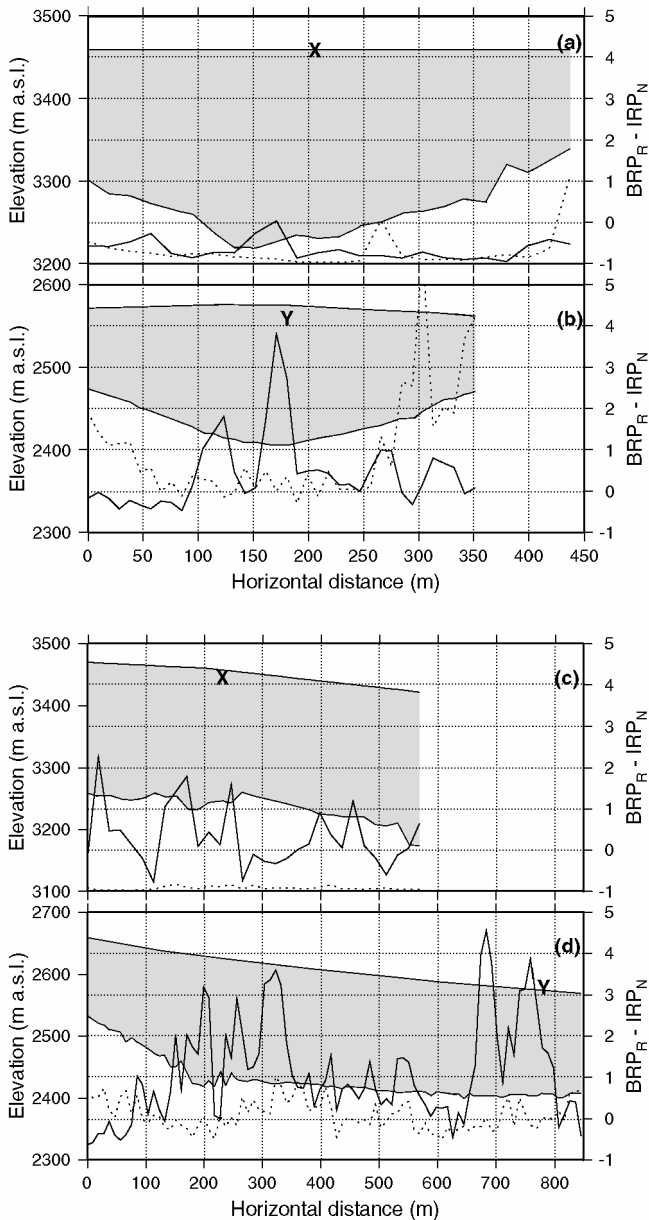


Fig. 7. Bedrock and surface along the two transverse RES profiles on Sofiyskiy glacier in the accumulation area (a) and the ablation area (b), and along two longitudinal RES profiles in the accumulation area (c) and the ablation area (d). The cross-profiles run from south to north; the long profiles run from west to east. X and Y denote the position of the longitudinal profiles and are shown in Figure 2. Also given are BRP_R (thin solid line) and IRP_N (dotted line).

antenna characteristics and surface coupling constant for all measurements, the only controls on the BRP are ice thickness, the loss due to two-way propagation through the ice (attenuation) and the strength of the basal reflection (Copland and Sharp, 2001). To determine the actual BRP variations due to variations in bed properties, we thus need to account for the effect of varying ice thickness. This is done by establishing an empirical relation between the bed reflection and the two-way travel time, using all traces of the four profiles (Fig. 8). To reduce the scatter due to differences in reflection properties of different materials, data over a homogeneous section in both the accumulation and ablation areas were selected. Data from the accumulation area were subsequently corrected for density variations in firn. Best-fit curves for both areas are also drawn in Figure

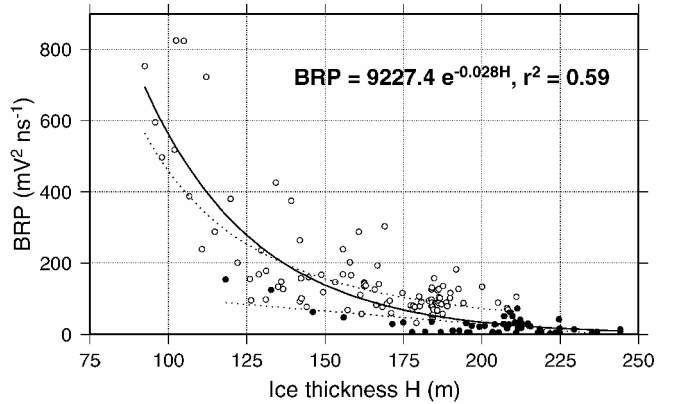


Fig. 8. Empirical relation between the bed reflection and the two-way travel time, using traces of the four profiles on Sofiyskiy glacier from the accumulation (filled circles) and the ablation area (open circles). Best-fit curves for both areas are given by dotted lines. The solid line displays the best-fit curve for all data. Values of the accumulation area were corrected for density variations in the firn to determine the ice thickness.

8. The best fit on all data was found to be an exponential function with $r^2 = 0.59$. Even though the accumulation data are corrected for density variations, the bulk of these data lie below the best-fit function, confirming a relatively low BRP for these measurements. Following Gades and others (2000), we now define the depth-corrected reduced BRP (BRP_R) as the measured BRP divided by the calculated BRP (according to the empirical relation displayed in Figure 8). Similarly, the normalized internal reflection power (IRP_N) is defined as the measured IRP divided by the mean IRP of all profiles. Although the BRP shows a high scatter, and the correlation coefficient of the best fit in Figure 8 is relatively low, the scaling operation has removed the major trend, i.e. the bias of high BRP associated with shallow ice thickness. The scatter will not hamper the interpretation of BRP_R , as we will only consider the relative importance of BRP_R in all four profiles.

In the accumulation area, values for BRP_R and IRP_N are distinctly lower than in the ablation area, and one can argue whether these differences are related to differences in englacial conditions (attenuation) or to differences in basal conditions (reflection). RES measurements on John Evans Glacier, Canada, also demonstrate this marked difference in BRP_R between accumulation and ablation area (Copland and Sharp, 2001). If englacial conditions are responsible for this difference, warm ice should correspond to low BRP_R . This contradicts the force-balance calculations which indicate that the accumulation area, where the lowest BRP_R is recorded, consists of stiff (cold) ice. It is therefore likely that BRP_R patterns are related to changes in basal conditions.

IRP_N variations seem to be primarily caused by differences in the snow and firn density, the ice temperature and the presence of occasional water inclusions and ice lenses. Since internal structures affect the emitted RES signal before it reaches the glacier bed, there is often an inverse relation between IRP_N and BRP_R . As such a relationship is not clear from our records, and IRP_N values are generally low for all profiles, the BRP_R appears little affected by internal attenuation processes. In the ablation area, two zones with a remarkably high basal reflection emerge from the RES profile. The first is located between 150 and 350 m (Fig. 7d) and coincides with the sudden decline of the bed

slope and the confluence area with side basins A and C. A number of hyperbolic structures can be identified which make it difficult to accurately determine the ice thickness. This could indicate a complex basal topography in this area and might influence the BRP_R as well. The second high basal reflection is located between 650 and 800 m (Fig. 7d) and might be related to the presence of several longitudinal crevasses and moulins on the glacier surface, enhancing surface water transport to the bed. The hyperbolic reflection structure near the glacier bed between 560 and 700 m (Fig. 6) may be related to this water transport as well. Since between 300 and 850 m along the longitudinal profile ice thickness, bed slope and englacial properties are relatively constant, the spatial variations in BRP_R seem to be caused by changes in the nature of the ice/bed interface. A striking feature in the ablation area is a strong BRP_R peak in the center of the transverse profile. It coincides with the downstream zone of high BRP_R in the longitudinal profile and is situated topographically at the valley bottom.

BASAL CONDITIONS

To allow for a more thorough assessment of basal conditions, we applied a three-layer reflectivity model (Born and Wolf, 1986) to calculate theoretical bed reflectivity values for ice, a basal layer and the bedrock floor, respectively. All layers were considered to be non-magnetic, homogeneous and smooth. The theoretical bed reflectivity \mathcal{R} of this three-layer case depends on the thickness h and the conductivity σ of the middle layer, as well as on the dielectric permittivity ε of each layer (see Appendix). The top layer has a value of $\varepsilon_1 = 3.17$ for temperate ice. The occurrence of gneiss, green schist and sandstone in the Sofiyskiy valley suggests a permittivity of $\varepsilon_3 = 9.0$ for bedrock. We performed model runs for a mixture of bedrock material and water with a varying porosity Φ . According to Looyenga (1965), the bulk dielectric permittivity is

$$\bar{\varepsilon} = \left[\Phi \varepsilon_w^{1/3} + (1 - \Phi) \varepsilon^{1/3} \right]^3, \quad (3)$$

where $\varepsilon_w = 88.0$ is the permittivity for water. For a high porosity ($\Phi \approx 1 \rightarrow \bar{\varepsilon} \approx \varepsilon_w$), Equation (3) represents the permittivity of a water layer, while for a low porosity ($\Phi \approx 0$), the permittivity is that of ice frozen to the bedrock. For $0 < \Phi < 1$, the bulk dielectric permittivity equals that of a till layer with varying water saturation. The bulk conductivity is also taken as a function of the porosity (Gades and others, 2000). For a given porosity Φ , \mathcal{R} will reach a maximum corresponding to a certain layer thickness h (Fig. 9). Since water has a very high dielectric constant relative to ice and rock, high values of BRP_R are consistent with the presence of a water layer, given the englacial conditions as discussed above. Besides a high reflective bed, high values for BRP_R could be due to an overestimated rate of attenuation with depth, which is unlikely as we notice high spatial variability in BRP_R for constant ice thickness, or could be due to multiple reflection, which only occurs in overdeepened areas. According to the model, the maximum \mathcal{R}_{max} occurs for a water layer with thickness $h = 1.5\text{m}$. However, for thinner and/or thicker water layers, a lower \mathcal{R} occurs. It is thus quite difficult to distinguish between a thin water layer and a saturated till layer. Low values of BRP_R are consistent with ice frozen to the bed, defined by a low dielectric constant.

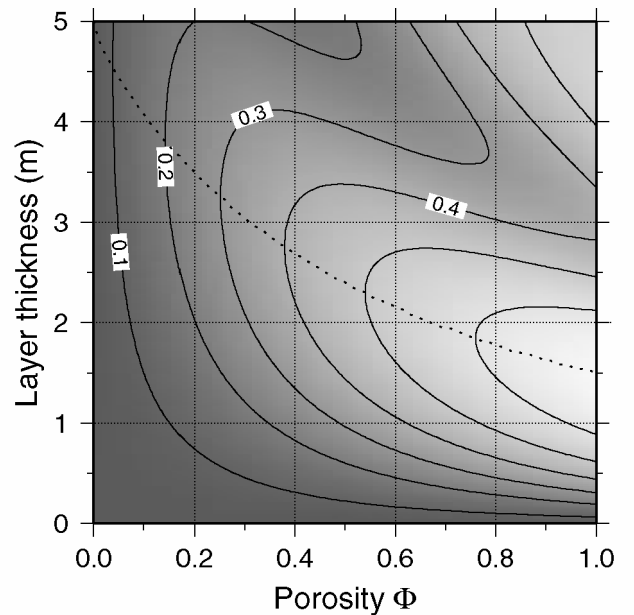


Fig. 9. Theoretical reflectivity \mathcal{R} as a function of porosity and layer thickness. The dotted line shows the path of the maximum reflectivity \mathcal{R}_{max} .

Water forms an important component of the glacier system in the ablation area during summer. Not only is surface melting a prominent feature, the glacier surface is characterized by a large number of surface meltwater channels and moulins; the latter are abundant in the downstream part of the ice tongue. At the front of the glacier, several meltwater rivers emerge from underneath the glacier ice. It is therefore likely that a well-developed subglacial drainage system exists. However, the sparsity of ice-thickness measurements outside the central flowline area do not allow for a reconstruction of the subglacial water network, based on the subglacial hydraulic potential (e.g. Copland and Sharp, 2001). Only the distinct peak in BRP_R in the cross-sectional profile of the ablation area, situated at the valley bottom (Fig. 7b), indicates that a subglacial water channel might be present. The presence of water — identified by spatially consistent high-reflectivity zones in the RES profiles — in the ablation area is in agreement with the force-balance calculations, which indicate that significant basal sliding should occur. Values of BRP_R in the accumulation area are, in view of the above discussion on englacial conditions (attenuation) and the generally low IRP_N , consistent with the force-balance analysis as well, which indicates that basal sliding does not occur. According to the reflectivity model, low BRP_R values point to a bedrock reflection, although thin layers of saturated till or water exhibit low reflections as well. Nevertheless, high basal reflectors are not encountered in the accumulation area, and since englacial conditions seem constant for both RES profiles, the presence of subglacial water can be excluded.

The marked difference in basal conditions between the accumulation and the ablation area, confirmed by the RES measurements and the force-balance analysis, might indicate that Sofiyskiy glacier exhibits a polythermal regime. Although, for a polythermal glacier, the ice is relatively shallow in the accumulation area ($< 250\text{m}$), the short summer season and extreme cold winter conditions, typical of a continental climate, might be held responsible for these conditions.

CONCLUSIONS

A 4 year field program carried out on Sofiyskiy glacier led to a proper determination of the physical characteristics of this so-called summer-accumulation-type glacier. Like many glaciers elsewhere on Earth, glaciers in this area have retreated significantly over the last century. We calculated a mean retreat rate of 18.3 m a^{-1} for Sofiyskiy glacier over a 100 year period. Sofiyskiy is at present characterized by a relatively low accumulation, the ELA being situated at approximately 3200 m a.s.l. Its surface velocity reaches a maximum near the icefall, about 110 m a^{-1} . A force-balance analysis indicates that significant basal sliding should occur in the ablation area (downstream of the icefall), while upstream of the icefall deformational velocities are consistent with stake measurements.

RES along several profiles on the glacier not only revealed the subglacial topography, but allowed for an analysis of basal characteristics through the BRP measured from the radar sounding data. The spatial patterns of high BRP likely indicate the presence of subglacial water in the ablation area. Low basal reflections dominate the accumulation area, which reflects a lack of subglacial water. Sofiyskiy glacier may therefore exhibit a polythermal regime.

ACKNOWLEDGEMENTS

This paper forms a contribution to the research project “Glacier Fluctuations and Climatic Change in South Siberia” (Federal Office for Scientific, Technical and Cultural Affairs), contract IN/RU/006(1998–2000), and the research project “Climatic Change in the Altai Mountains” (Scientific Research Fund, Belgium (FWO)), contract G.0120.97 (1998–2000). The authors are indebted to I. Novikov, M. Mokhtar and M. Monté for assistance in the field, K. Matsuoka for critical remarks on a first draft of this paper, and R. C. A. Hindmarsh, A. Hubbard and F. Navarro for helpful comments during reviewing.

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APPENDIX

THREE-LAYER REFLECTIVITY MODEL

Consider a plane-parallel absorbing film situated between two dielectric media. The total reflectivity is then given by (Born and Wolf, 1986)

$$\mathcal{R} = |r|^2 = \left| \frac{\hat{\rho}_{12} e^{i\phi_{12}} + \hat{\rho}_{23} e^{-2v_2\eta} e^{i(\phi_{23} + 2u_2\eta)}}{1 + \hat{\rho}_{12}\hat{\rho}_{23} e^{-2v_2\eta} e^{i(\phi_{12} + \phi_{23} + 2u_2\eta)}} \right|^2, \quad (4)$$

where $\eta = 2\pi h/\lambda_0$, h is the layer thickness and $\lambda_0 = 60 \text{ m}$ is the wavelength of the radar signal in free space. $\hat{\rho}_{12}, \hat{\rho}_{23} \in \mathbb{C}$ are the ratios of the electric vectors, which are a function of the angle of incidence and the refractive index. Indices 1–3 refer to the layer numbers. For an electric vector perpendicular to the plane of incidence, the ratio of the electric vectors is written as

$$\hat{\rho}_{12} e^{i\phi_{12}} = \frac{n_1 \cos \theta_1 - (u_2 + iv_2)}{n_1 \cos \theta_1 + (u_2 + iv_2)}, \quad (5)$$

where $(u_2 + iv_2)^2 = \hat{n}_2^2 - n_1^2 \sin^2 \theta_1$ and θ is the angle of incidence. n is the refractive index, which for a non-conducting medium is defined as $n = \sqrt{\mu\epsilon}$, where μ is the magnetic permeability and ϵ the dielectric permittivity. For a conducting material, the refractive index \hat{n} is a function of the conductivity as well, through the attenuation index κ (Born and Wolf, 1986)

$$\hat{n}_2^2 = n_2^2 (1 + 2i\kappa_2 - \kappa_2^2) = \mu_2 \hat{\epsilon}_2 = \mu_2 \left(\epsilon_2 + i \frac{\mu_2 \sigma}{\nu} \right), \quad (6)$$

where σ is the dielectric conductivity and ν the frequency. In Equation (4), ϕ is the phase change defined by

$$\tan \phi_{12} = \frac{2v_2 n_1 \cos \theta_1}{u_2^2 + v_2^2 - n_1^2 \cos^2 \theta_1}. \quad (7)$$